

A  $\sim 15$  kpc outflow cone piercing through the  
halo of the blue compact dwarf galaxy  
SBS 0335-052E

Escape of Lyman radiation from galactic labyrinths  
(Crete, April 18-21, 2023)

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April 20, 2023

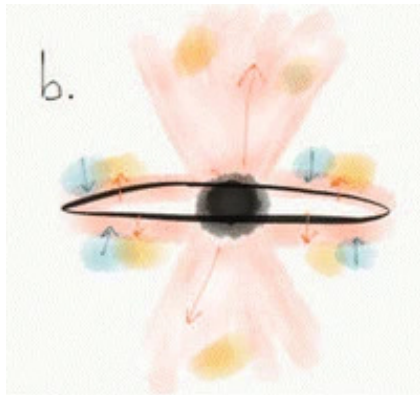


# Outflows: Key process for understanding galaxies.

Star-formation driven feedback and outflows regulate availability and chemical composition of gas for forming stars (in concert with inflows).

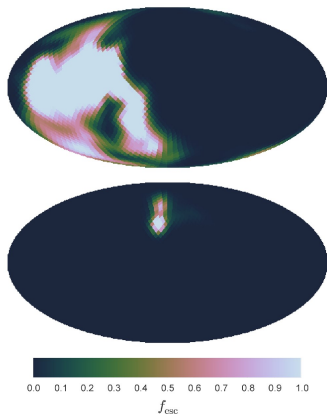
## Impact

- star-formation history
- size
- colour
- metallicity
- inner dark-matter density profiles
- ...



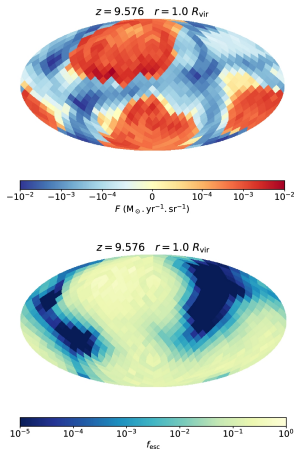
# Outflows are important in regulating $f_{\text{esc}}^{\text{LyC}}$ & $f_{\text{esc}}^{\text{Ly}\alpha}$

Beamed outflows: LyC/Ly $\alpha$  escape into small solid angles.



2 massive haloes  
( $M_H \sim 10^7 M_\odot$ ) at  $z \sim 12$

Paardekooper et al. (2015)

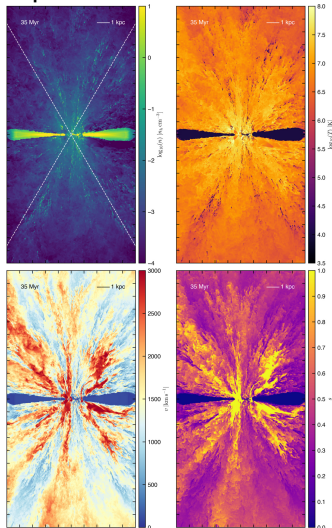
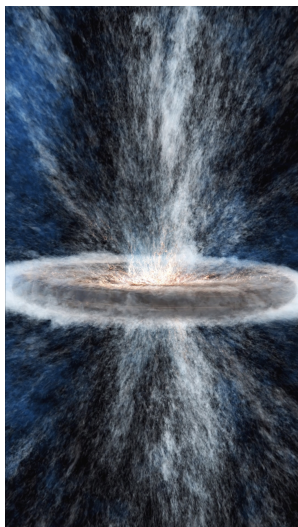


$M_H \sim 10^8 M_\odot$

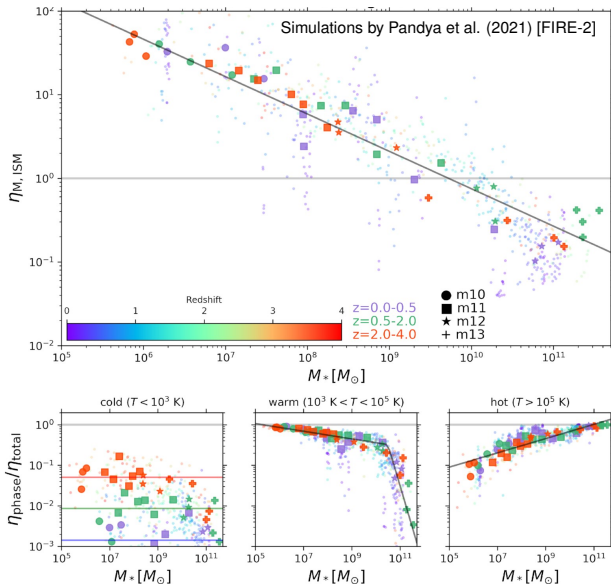
Trebitsch et al. (2017)

# Outflow simulation of an individual galaxy

GPU based - run on Titan super computer - assumes  $\alpha$



Key quantity to characterise winds:  $\eta = \dot{M}_{\text{wind}} / \dot{M}_{\text{SFR}}$ .

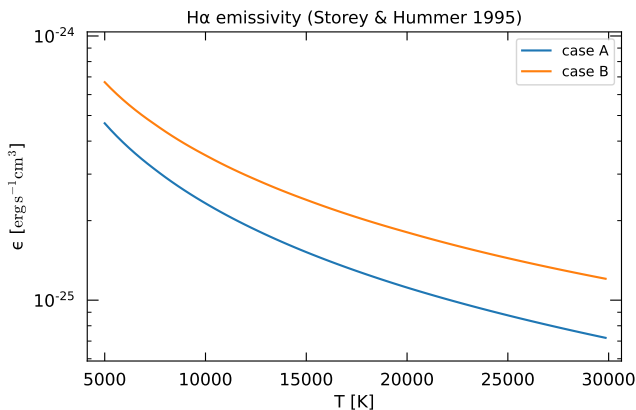


Starburst galaxies: most mass loaded in the warm phase.

# Mapping the diffuse ionised gas - H $\alpha$

H $\alpha$  ( $\lambda 6563$ ) from volume  $V$ , density  $n$  ( $\sim n_e \sim n_p$ ), temperature  $T$  in recombination equilibrium ( $E_{\text{ion}} = 1 \text{ Ry}$ ):

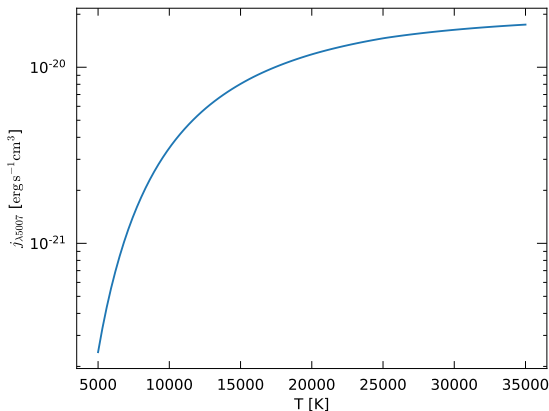
$$L_{\text{H}\alpha} = \epsilon_{\text{H}\alpha}(T) \times n^2 \times V \quad (1)$$



# Mapping the diffuse ionised gas - [o III] $\lambda 5007$

Collisional Excitation of  $O^{2+}$  ( $E_{\text{ion}} = 2.58 \text{ Ry}$ )  $\rightarrow$  “Forbidden” Line (magnetic dipole transition).

$$L_{\lambda 5007} = j_{\lambda 5007}(T) \times n_e \times n(O^{2+}) \times V \quad (2)$$

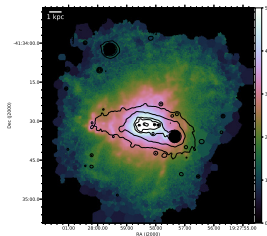


Collisional emissivity from `pyneb` (Luridiana, Morisset & Shaw 2015).

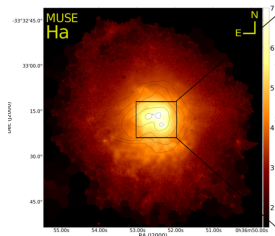


# VLT/MUSE – ideal instrument to map outflows spatially

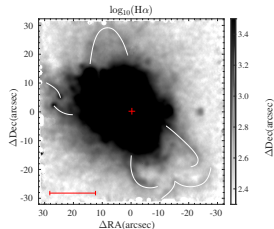
Unprecedented sensitivity for low-SB line emission.



ESO 338 -  $d = 40$  Mpc  
 $M_{\text{dyn}} \sim 10^8 M_{\odot}$   
(Bik et al. 2018)



Haro 11 -  $d = 87$  Mpc  
 $M_{\text{dyn}} \sim 10^{10...11} M_{\odot}$   
(Menacho et al. 2019)

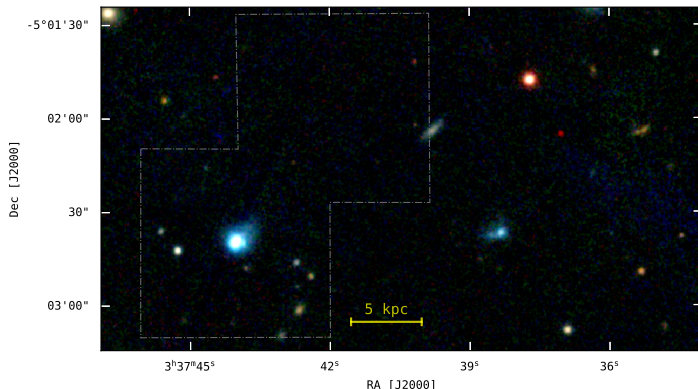


Haro 14 -  $d = 13$  Mpc  
 $M_{\text{dyn}} \sim 10^8 M_{\odot}$   
(Cáirois et al. 2022)

Seems like every star-bursting dwarf observed with MUSE shows extended H $\alpha$  structures (relative to continuum).

# The SBS 0335-052 system ( $d = 58 \text{ Mpc} / z = 0.0135$ )

A “reference laboratory” for understanding high- $z$  galaxies

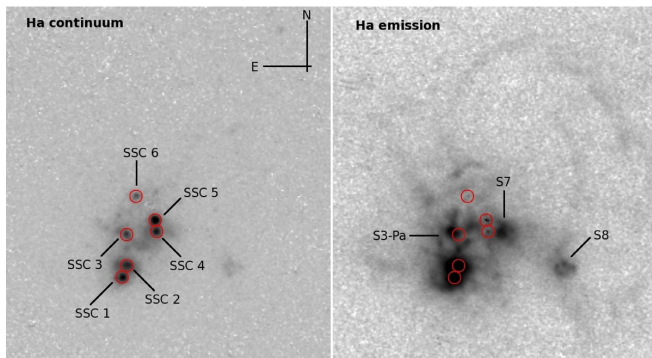


SBS 0335-052E:  $M_{UV} = -16.4$ ,  $M_{\star} = 5.7 \times 10^6 M_{\odot}$ ,  
 $SFR = 1.2 M_{\odot} \text{ yr}^{-1}$ ,  $12 + \log(\text{O}/\text{H}) = 7.25$

Izotov et al. 1990, Nature 343, 238.

(SBS 0335-052W:  $12 + \log(\text{O}/\text{H}) = 6.86 \dots 7.22$ )

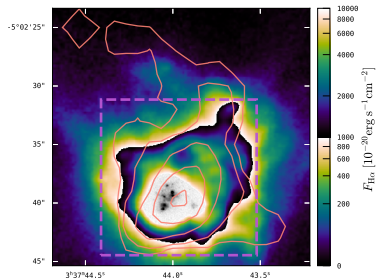
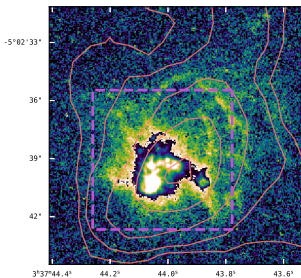
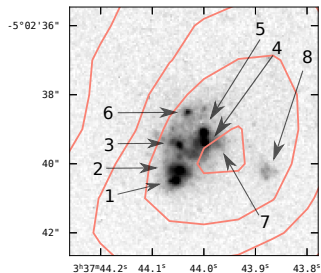
# HST: Age Gradient (SE→NW) & H $\alpha$ super-shell.



Properties of the Six SSCs. SED at Fixed Age for Mass Determination. Age is Determined Through EW(H $\alpha$ )

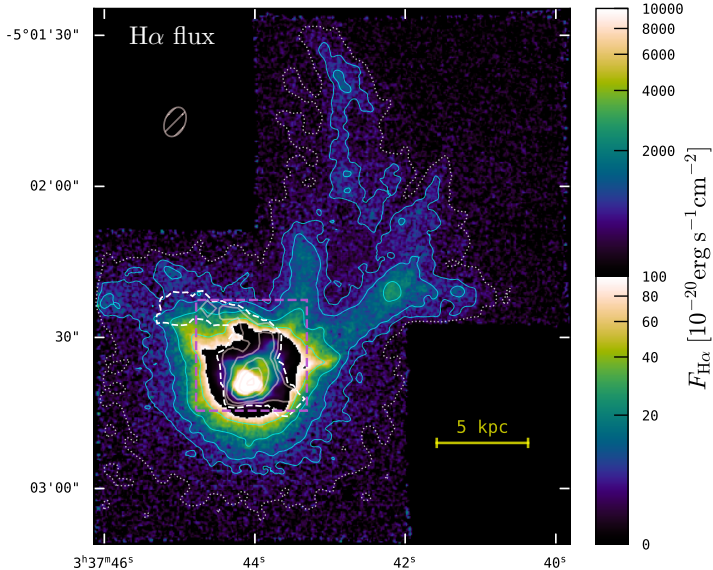
IDs	$f(\text{H}\alpha)$ ( $\text{erg s}^{-1} \text{cm}^{-2}$ )	EW (H $\alpha$ ) ( $\text{\AA}$ )	Age (R08) <sup>a</sup> (Myr)	Age <sup>b</sup> (Myr)	Mass ( $M_{\odot}$ )	$A_V$ (mag)	$R_{\text{HII}}$ (pc)
SSC1	$5.14 \times 10^{-14}$ (1%)	3100.0(134.4)	$\leq 3.3$	3.0	$4.7 \times 10^5$	0.73	11.3(11.5) <sup>c</sup>
SSC2	$3.34 \times 10^{-14}$ (1%)	2300.0(104.7)	$\leq 3.4$	3.0	$3.7 \times 10^5$	0.65	10.4(9.7) <sup>c</sup>
SSC3	$5.77 \times 10^{-15}$ (2.5%)	787.0(53.4)	6.8(2.5)	7.0	$7.1 \times 10^5$	0.92	29.1(26.6) <sup>d</sup>
SSC4	$2.42 \times 10^{-15}$ (3.9%)	176.0(10.7)	12.4(1.7)	11.0	$1.1 \times 10^6$	0.20	32.9(26.9) <sup>e</sup>
SSC5	$2.29 \times 10^{-15}$ (4.5%)	84(4.4)	15.1(2.3)	13.0	$2.9 \times 10^6$	0.84	38.6(29.5) <sup>e</sup>
SSC6	$6.34 \times 10^{-16}$ (7.7%)	187(22.5)	13.9(1.9)	11.0	$2.6 \times 10^5$	1.08	20.6(21.49) <sup>e</sup>

# Let's zoom out (HST FR565N, F550M, & now MUSE).



$$t_{\text{exp}} \approx 1.5 \text{ h}$$

$$N_{\text{H}\alpha} = \{5, 10, 20, 30, 40\} \times 10^{20} \text{ cm}^{-2}$$

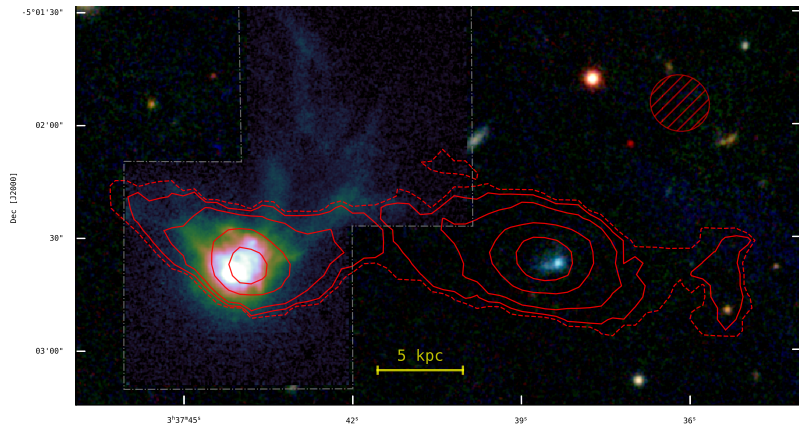


$1'35'' \times 1'46''$  (26.2 kpc  $\times$  29.2 kpc).

Contours:  $SB_{H\alpha} = \{0.75, 1.5, 2.5, 5, 12.5\} \times 10^{-18} \text{ erg s}^{-1} \text{ cm}^{-2} \text{ arcsec}^{-2}$  (dotted white & cyan),

$N_{HI} = \{2.5, 5, 10, 20, 30, 40\} \times 10^{20} \text{ cm}^{-2}$  (dashed white & grey).

# Structure extends outward of HI Halo

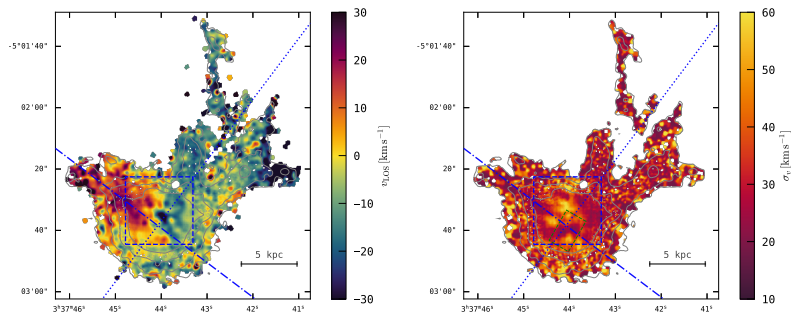


$H\alpha$  – log-stretch from 0 to  $1.25 \times 10^{-15} \text{ erg s}^{-1} \text{ cm}^{-2}$ .

21 cm – contours  $N_{\text{HI}} = \{0.5(2.5\sigma), 1, 2, 5, 8\} \times 10^{20} \text{ cm}^{-2}$ .

Background: grz colour composite (Pan-STARRS).

# Kinematics of H II

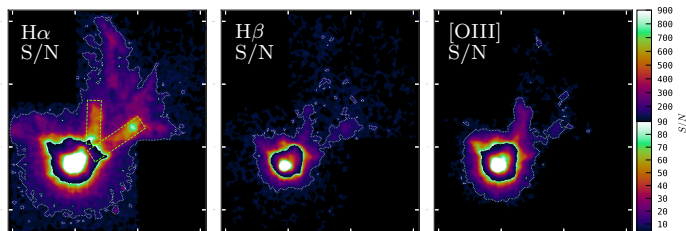


Filaments characterised by:

- Narrow lines –  $\sigma_v = 20 \dots 30 \text{ km s}^{-1}$ .
- No velocity gradient in outwards direction. Perhaps small  $\Delta v \sim 5 \text{ km s}^{-1}$  between them.

Symmetric bifurcation wrt. to minor axis. *Clue for outflow.*

# MUSE Detections of $H\alpha$ , $H\beta$ , $[O III] \lambda 5007$

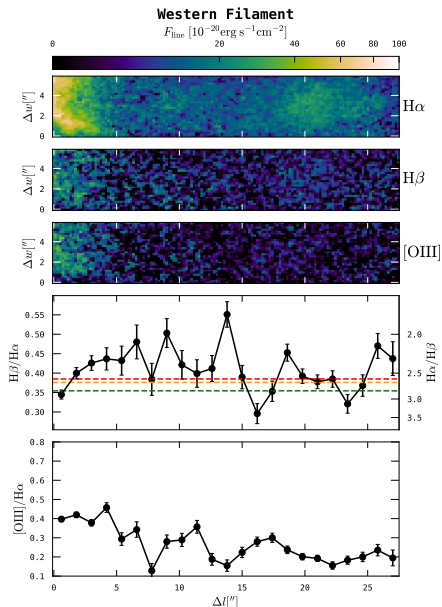
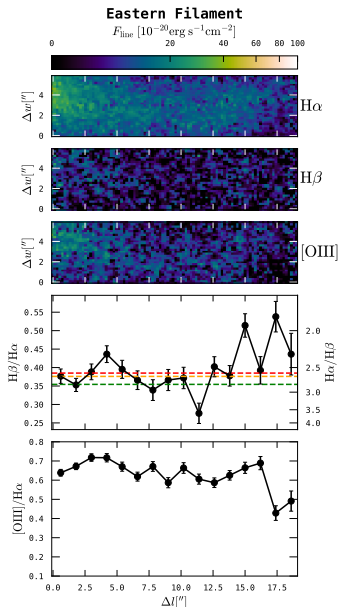


Contour:  $SN_{LSD} = 8$ .

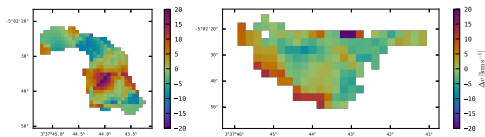
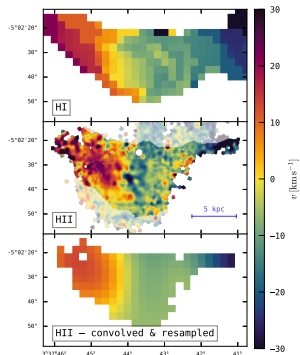
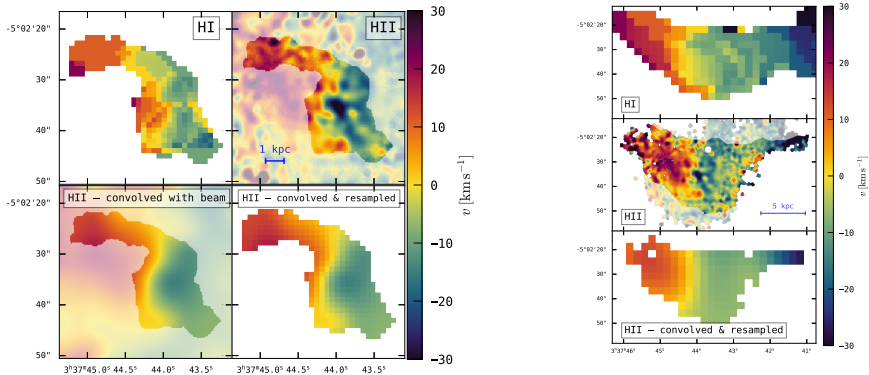
S/N after 3D Gaussian filter ( $\sigma_v = 150 \text{ km s}^{-1}$ ,  $\sigma_{2D} = 2''$ )  
LSDCat (Herenz, E. & Wisotzki L. 2017),  
see also LSDCat2.0 (Herenz 2023, see poster).



# Line Ratio Analysis — $H\beta/H\alpha$ & $[O\text{ III}]/H\alpha$



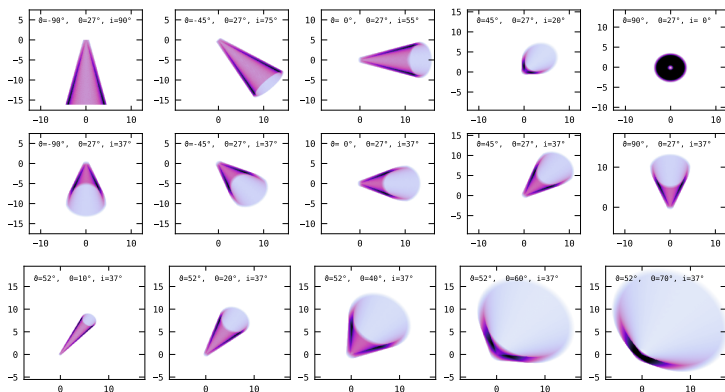
# HI vs. HII kinematics



Bulk of H II not comoving with H I.

# What are the filaments? Cone-Wall Structure.

Geometrical Parameters:  $\theta$  (opening angle),  $\vartheta$  (position angle),  $i$  (inclination),  $h$  (height),  $t$  (wall thickness).



Inside:  $V_{\text{hot}} = \pi/3 \cdot \tan^2(\theta/2) \cdot h^3$

Walls:  $V_{\text{HII}} = \pi/3 \cdot \tan^2(\theta/2) \cdot ((h + \Delta h)^3 - h^3 - \Delta h^3)$ , with  $\Delta h = 2 \cdot t \cdot \cos(\theta/2)$ .

# Volumes $V_{\text{HII}}$ and $V_{\text{hot}}$ .

Fix geometrical parameters from observations:

- $\vartheta = 52^\circ$ ,  $i = 43^\circ$  (Moiseev et al. 2010).
- $\theta = 2 \times \arctan[\sin(i) \times \tan(\theta_P/2)] = 27^\circ$ , where  $\theta_P = 34^\circ$  is the observed opening angle.
- $h = l_p \cdot \cos(\theta/2) / \sin(i) = 16.2 \text{ kpc}$ , where  $l_p = 10 \text{ kpc}$  is the projected length of the filaments.
- $t = 1.5 \text{ kpc}$ .

$$\underline{V_{\text{hot}} = 256 \text{ kpc}^3} \quad \underline{V_{\text{HII}} = 142 \text{ kpc}^3} \quad (3)$$

# Hot phase: Mass, Mass Loading ( $\eta$ ), & $v_{\text{wind}}$

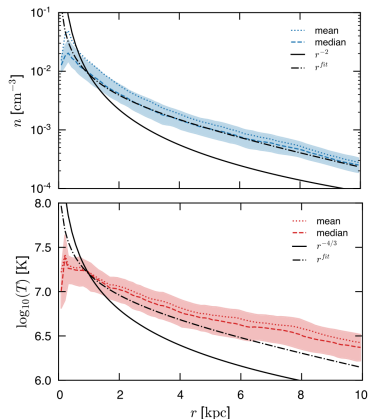
$$M_{\text{hot}}(V_{\text{hot}}, n) = 6 \times 10^{5 \dots 6} M_{\odot}$$

$$\underline{\eta_{\text{hot}} \lesssim 0.1}$$

in  $M_{\text{hot}} = \int \eta_{\text{hot}} \dot{M}_{\star} dt$  with  $\Delta t = 10 \text{ Myr}$   
and  $\dot{M}_{\star} \approx 1 M_{\odot} \text{ yr}^{-1}$

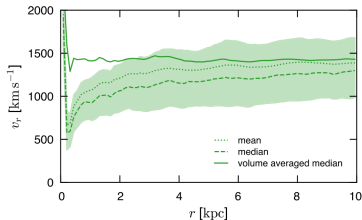
$$\begin{aligned} v_{\text{wind}} &= h / (t_{\star} - t_{\text{SN}}) \\ &= 16.2 \text{ kpc} / 10 \text{ Myr} \\ &= 1620 \text{ km s}^{-1} \end{aligned}$$

for  $t_{\star} = 15 \text{ Myr}$  and  $t_{\text{SN}} = 5 \text{ Myr}$ .



$$\begin{aligned} T_{\text{hot}} &\sim 10^6 \dots 10^7 \text{ K}, \\ n &\sim 10^{-3} \dots 10^{-4} \text{ cm}^{-3}. \end{aligned}$$

(CGOLS; Schneider et al. 2020)



# Warm phase: Mass, Mass Loading, Observability

Fix  $n$  and  $T$  in Eq. (1), such that we reproduce the observed flux  $F_{\text{H}\alpha}$  with  $V_{\text{HII}}$  from Eq. (3) (assuming Case-A):

$$n = \{3, 4, 5\} \times 10^{-2} \text{ cm}^{-3} \quad \text{for} \quad T = \{1, 1.5, 2\} \times 10^4 \text{ K} \quad (4)$$

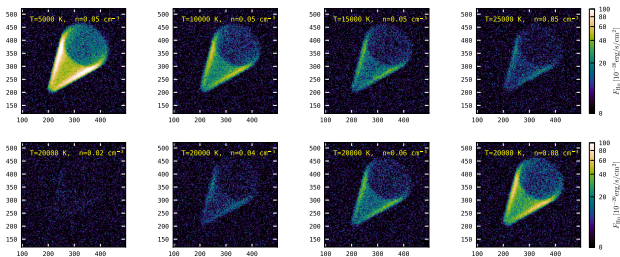


Figure: Simulated MUSE H $\alpha$  observations of cone-wall structure for varying  $T$  and  $n$ .

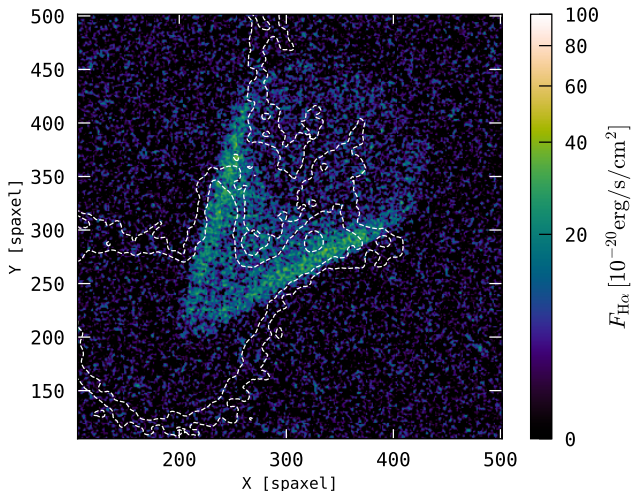
Best match to observations:

$$n = 5 \times 10^{-2} \text{ cm}^{-3} \quad \& \quad T = \{1, 1.5, 2\} \times 10^4 \text{ K}$$

$$\Rightarrow \underline{M_{\text{HII}} = 1.5 - 1.7 \times 10^8 M_{\odot}} \quad \underline{\eta_{\text{HII}} \gtrsim 10}$$

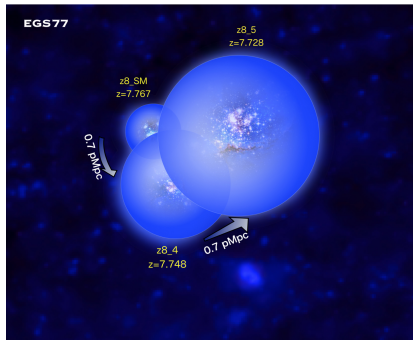
# Toy model consistent with observations and theory

Provides realistic masses and loading factors. Required  $T$  and  $n$  consistent with expectations for “hot” and “cool” phase of outflow. Reproduces basic H II morphology.



# Implications for the Epoch of Reionisation?

Reionisation often conceptualised via “ionised” bubbles.

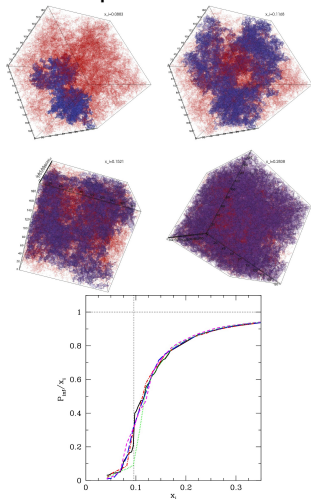


Tilvi et al. (2020, ApJL 891, L10): 3 Ly $\alpha$  emitting galaxies at  $z = 7.7$  (Keck/MOSFIRE)

$$R_S = \left( \frac{3\dot{N}_{\text{ion}} f_{\text{esc}} t}{4\pi n_{\text{HI}}(z)} \right)^{1/3}$$

Evolution of EoR: pre-overlap  
→ overlap → post-overlap...

EoR is phase transition.



Furlanetto & Oh (2016) - spherical symmetry in those models: “The overlap phase –  $0.1 \lesssim x_i \lesssim 0.9$  – sees the ionized and neutral gas almost entirely contained in just two distinct, delicately intertwined regions.”



## Summary on SBS0335-052E

- SBS0335-052E shows two low-SB  $H\alpha$  filaments ( $SB_{H\alpha} \approx 1.5 \dots 3 \times 10^{-18} \text{erg s}^{-1} \text{cm}^{-2} \text{arcsec}^{-2}$ ), that extend outside of H I halo.
- Balmer decrement indicative of low optical depth (neither Case-A nor Case-B).
- Direction and symmetry with respect to kinematic model of the galaxy: Clue for Outflow.
- Geometrical toy model reproduces observables, when fed with input parameters ( $n, T$ ) according to theoretical expectations for the wind.